



## **Drag Reduction of Hypersonic Spiked Blunt Body with Sideward Jets**

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### **Abstract**

Drag reduction for blunt body flying at hypersonic speed were numerically investigated in this paper, with spikes and sideward jets applied to re-configured the flow field. The spikes and the sideward jets led to the formation of the conical shock, the reattached shock, and the circumfluence, instead of a bow shock in front of the blunt body. The flow field reconfiguration played a dominant role in drag reduction. This paper proposed a parameter for the design of drag reduction of the spiked bodies, by combing the locations of the reattached shock and its interaction with the conical shock.

**Keywords:** *hypersonic, drag reduction, spike, sideward jet, numerical simulation*

### **Nomenclature (Tahoma 11 pt, bold)**

$Ma$  – Mach number of the incoming flow  
 $L$  – The length of the spike  
 $D$  – The diameter of the blunt body (sphere cylinder)

$y_1$  – The location of the reattached shock  
 $y_2$  – The location of the interaction between the reattached shock and the conical shock

### **1. Introduction**

The concept of a spiked blunt-body was first proposed by Bogdonoff et al. [1]. Investigations showed that spikes can construct a separation zone over blunt bodies to significantly lower the aerodynamic heat rate and pressure distribution, which was valuable for thermal protection and drag reduction. Han set up a new non-ablation concept utilizing adaptive drag reduction and thermal protection system (NADTPS) for hypersonic vehicles basing on a spiked blunt body with sideward jet nearby the spike nose [2]. This paper investigated the spiked blunt bodies with and without sideward jets flying at hypersonic speed, using numerical simulations. The incoming flow Mach number changed from 5 to 10, and the spike length ratio over the blunt body diameter changed from 0.5 to 1.0. The flow structure reconfiguration, especially the shock wave structure, was analyzed by flow field visualization. Aerodynamic force characteristics were calculated and compared with those of blunt bodies to analyze drag reduction effects and the corresponding mechanisms.

### **2. Governing Equations and Numerical Methods**

The two-dimensional axisymmetric compressible Navier-Stokes equation set was applied as the governing equations, with the thermally perfect gas model and ideal gas state equation. The 2<sup>nd</sup> order DCD scheme [3] was adopted to discrete the convective terms, and the Steger-Warming method [4] was applied for flux vector splitting. The 2<sup>nd</sup> order central difference scheme was used for viscous term discretization. The local time step method [5] was introduced to accelerate the simulation convergence process using different time steps at grid points.

### **3. Result and Discussion**

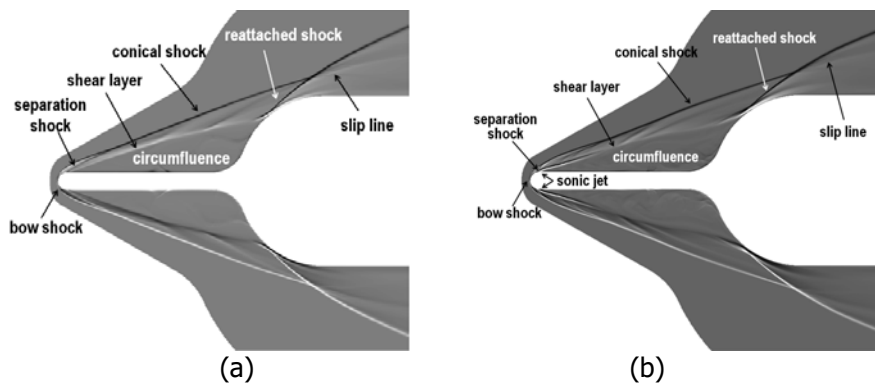
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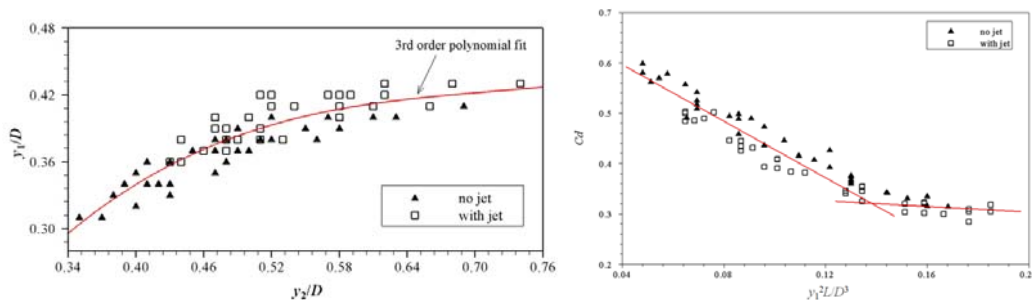
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The typical shock wave structures of pure spiked blunt bodies with defined labels are shown in Fig. 1(a) and Fig.1(b). Density gradient in the y direction was used to describe a schlieren photograph and re-configured flow field. The spike nose became the new stagnation with a little bow shock in front and conical shock around. The reattached shock was formed on the blunt-body shoulder, resulting in a zone with high pressure. The inverse pressure gradient caused some of the gas flow to the spike nose, in the opposite direction of incoming flow. Thus, a circumfluence was generated around the spike and the blunt-body nose.

The tendency of the 3<sup>rd</sup> order polynomial curve shows that there might exist a limited value of  $y_1/D$ , which maintained the reattached shock position as the spike length increased. Based on the numerical results, the limit value was about 0.43, located at the angle about 60-degree to the central line, as shown in Fig.2(a). Different from the reattached shock, the interaction between the conical shock and the reattached shock moved away from the central line as the spike length increased, so does the slip line resulted by the interaction. The drag coefficients in the 72 cases of the pure spiked bodies and the spiked bodies with sonic sideward jets were plotted together, as shown in Fig.2(b). With  $y_1^2L/D^3$  taken as the horizontal coordinate, the distribution of data points was significantly changed and more concentrated in the two drawn lines. Thus, the parameter  $y_1^2L/D^3$  might be a simple and suitable description for the two main factors during flow field reconfiguration. In addition, drag coefficient tendency changes and becomes more gently when  $y_1^2L/D^3$  is greater than 0.14. The reattached shock and the circumfluence criterion for drag reduction could be used by determining the critical value. Therefore,  $y_1^2L/D^3 \approx 0.14$  can be considered as a criterion for the design optimization regarding the spike length and the strength of sideward jets.



**Fig 1.** Typical flow field wave structure around the spiked blunt-body ( $Ma=6.0$ ,  $L/D=1.0$ ): (a) without sideward jet; (b) with sideward sonic jet nearby the spike nose.



**Fig 2.** The combination of the parameter for the design of drag reduction system: (a) A third-order polynomial fit  $y_1/D$  and  $y_2/D$ ; (b) Drag coefficients vary with the combined parameter  $y_1^2L/D^3$ .

**References**

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